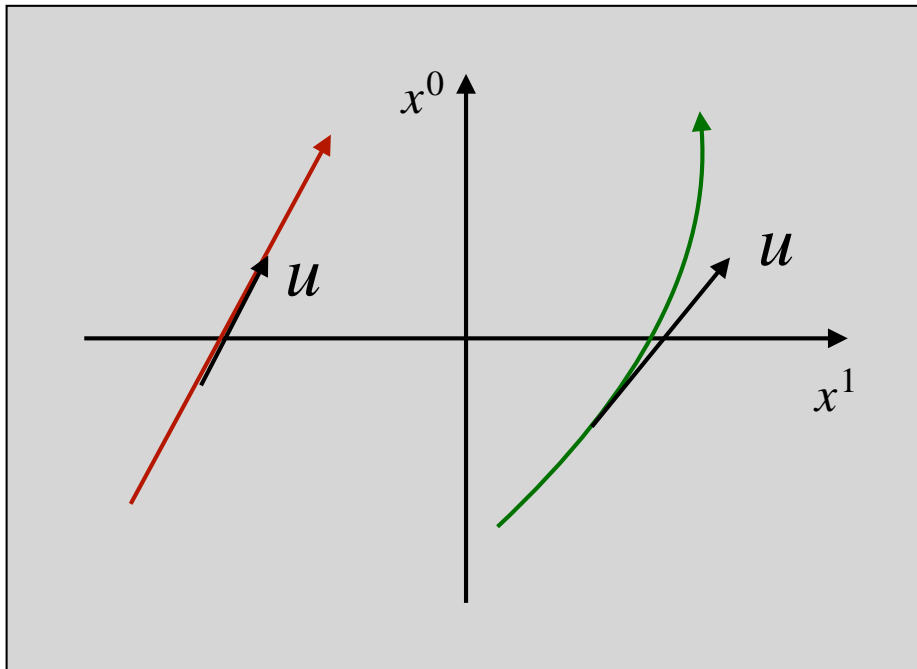


Kinematics in SR



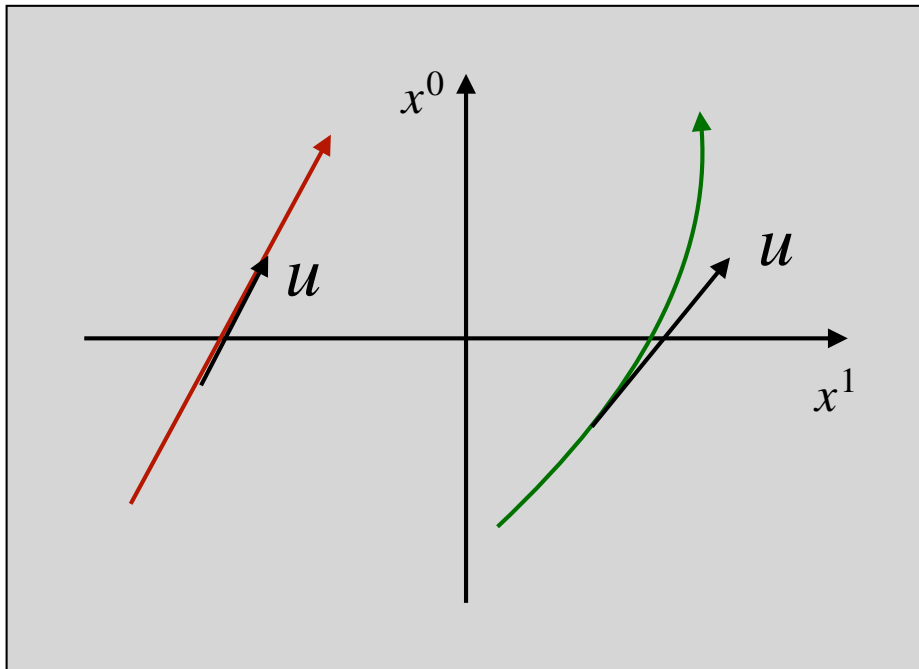
4-velocity u^μ

normalized $u \cdot u = -1$

future-pointing $u^0 > 0$

any co-moving frame has $e_0 = u$

Kinematics in SR



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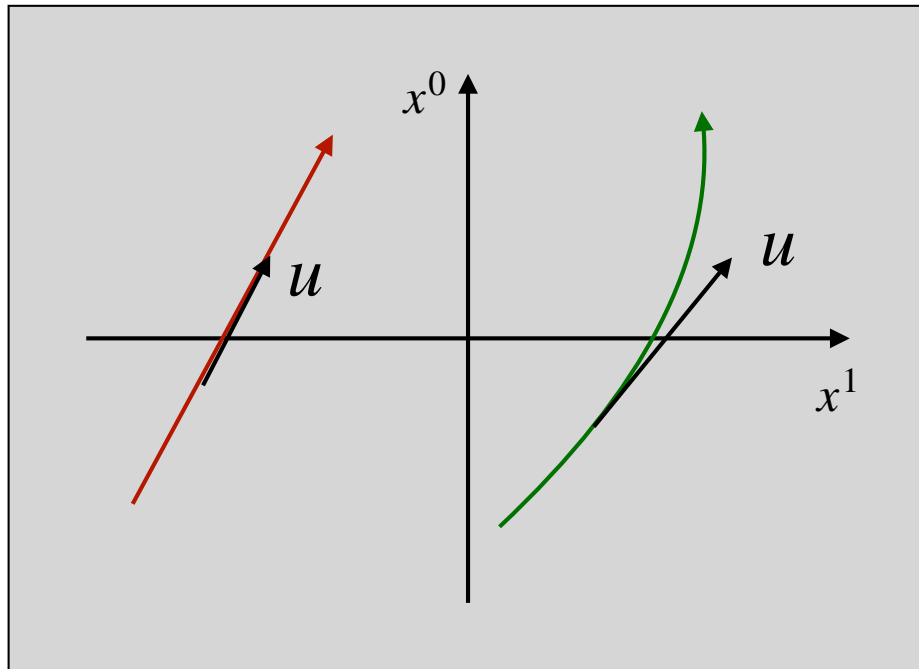
any co-moving frame has $e_0 = u$

proper time τ :

its flow matches the coordinate time of the co-moving frame

interpretation: time measured by a moving perfect clock

Kinematics in SR



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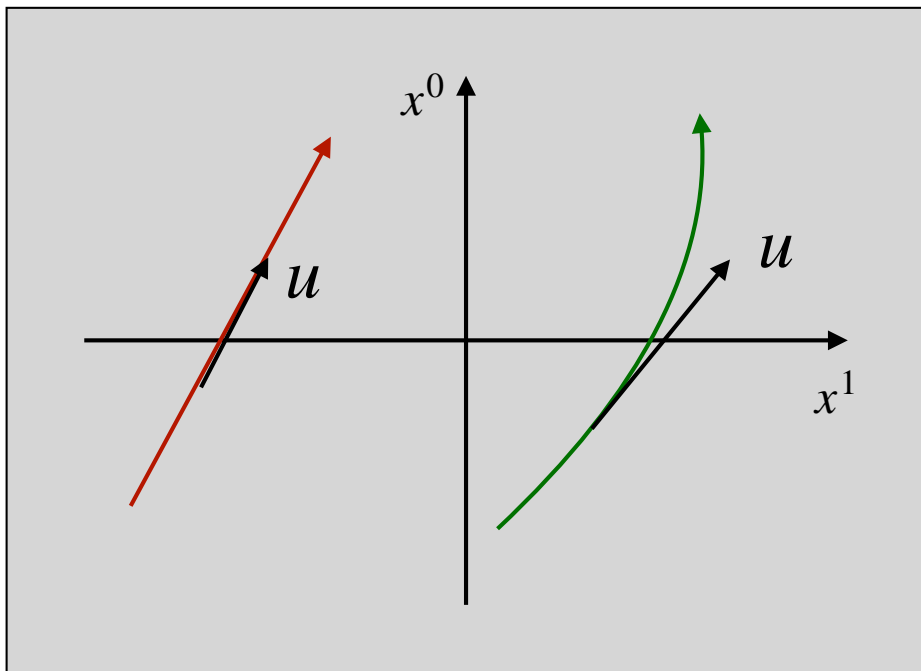
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$$x^0 = \gamma \tilde{x}^0 + \gamma \vec{v} \cdot \vec{\tilde{x}} \implies dx^0 = \gamma d\tilde{x}^0 = \gamma d\tau$$

Kinematics in SR



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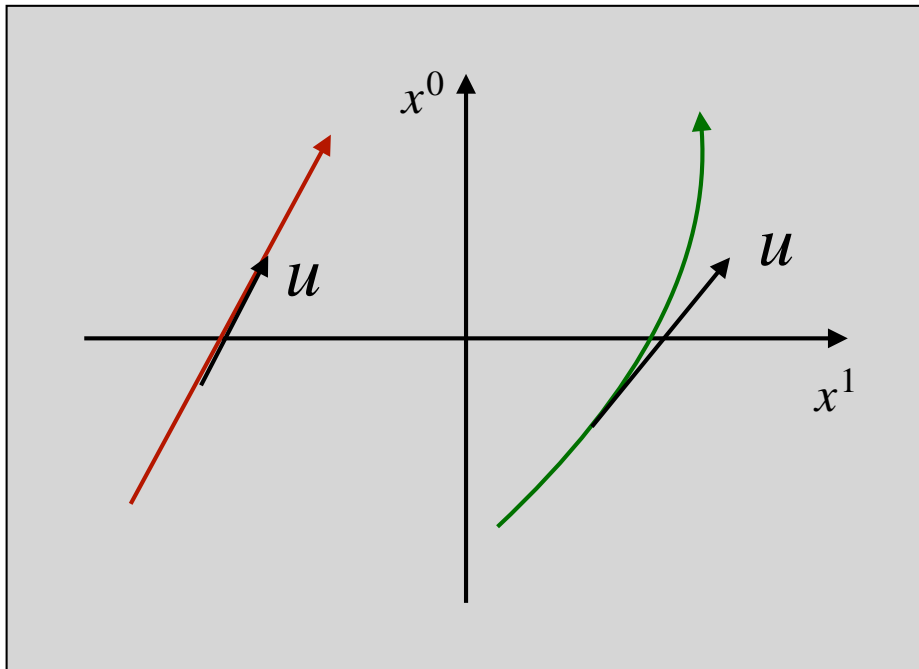
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Kinematics in SR



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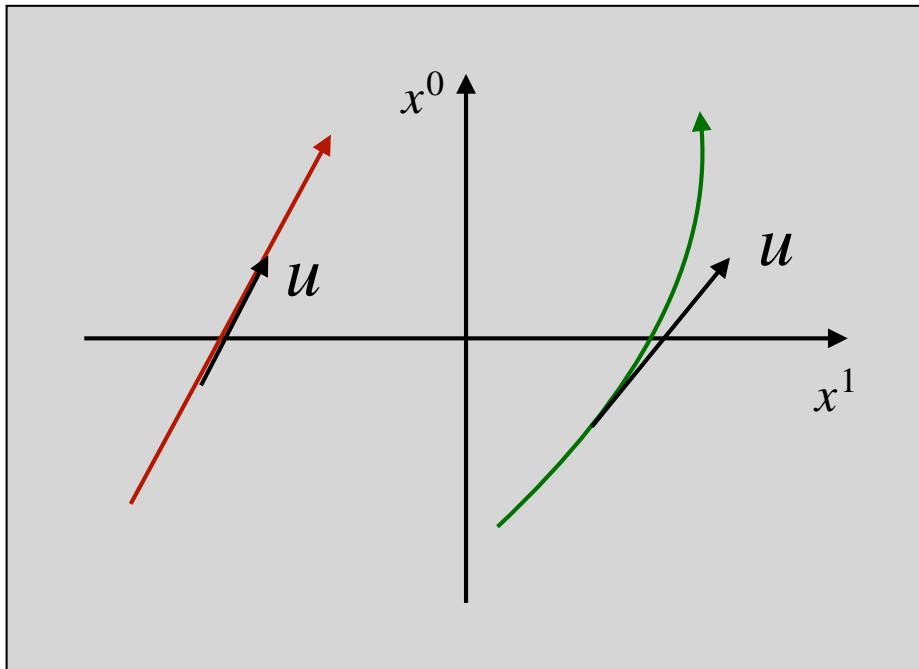
defined intrinsically, by the motion of the body

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Kinematics in SR



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defined intrinsically, by the motion of the body

natural parametrization related to 4-velocity

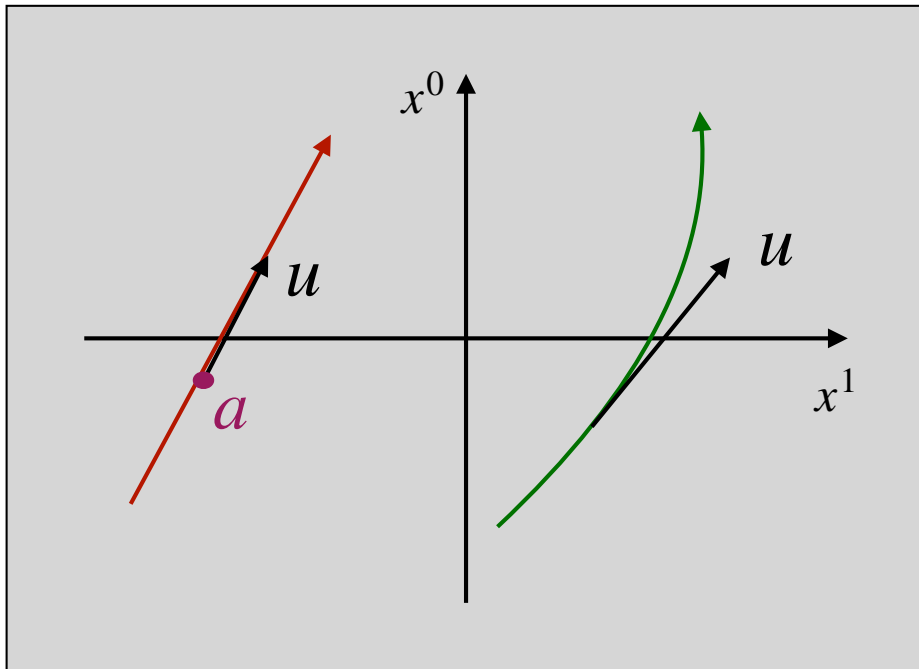
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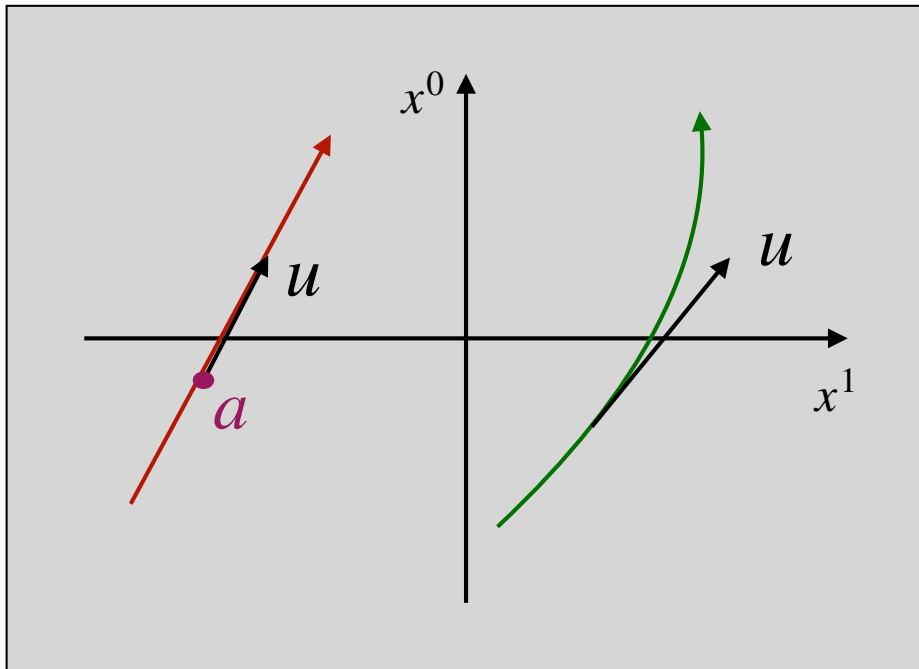
$$x^\mu(\tau) \quad \frac{dx^\mu}{d\tau} = \frac{dx^\mu}{dx^0} \frac{dx^0}{d\tau} = \begin{pmatrix} 1 \\ v^i \end{pmatrix} \gamma = u^\mu$$

Kinematics in SR



**massive particle moving with
constant velocity**

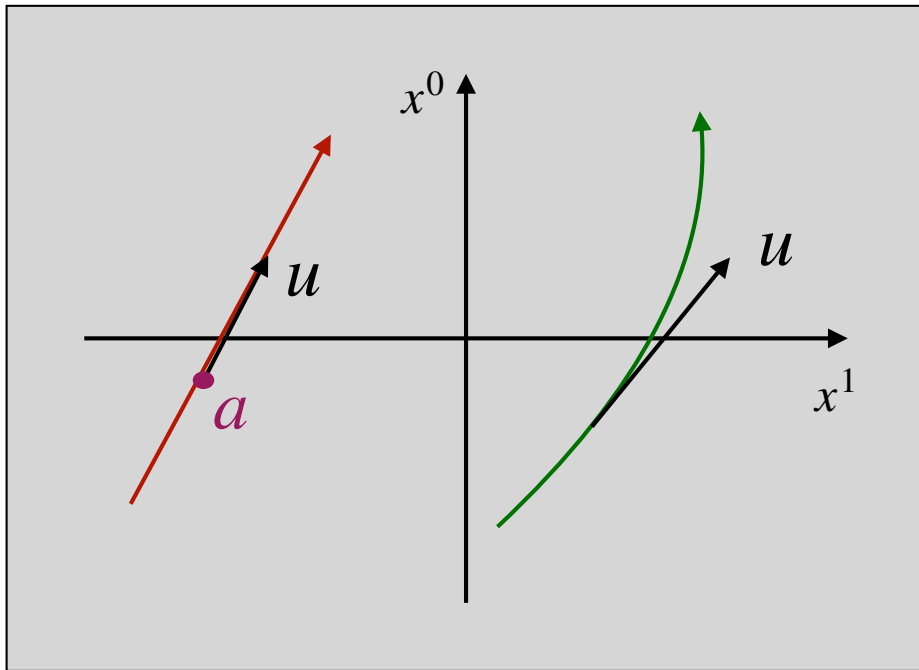
Kinematics in SR



**massive particle moving with
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$$x^\mu(\tau) = a^\mu + \tau u^\mu$$

Kinematics in SR



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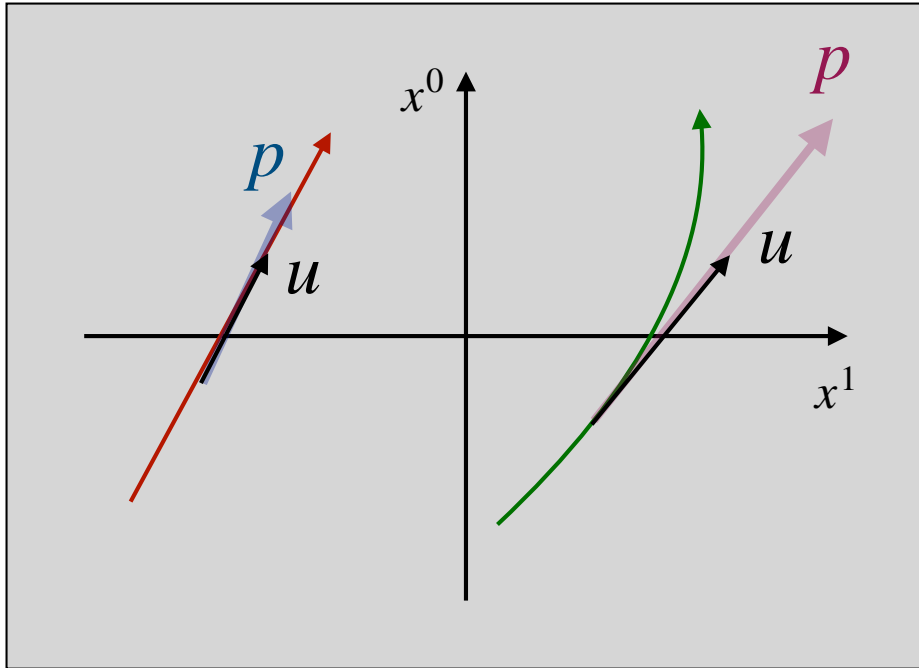
$$x^\mu(\tau) = a^\mu + \tau u^\mu$$

reparametrization:

$$\tau \rightarrow \tilde{\tau} = \tau + D$$

$$a \rightarrow \tilde{a} = a - D u$$

Kinematics in SR



4-momentum

particle of rest mass m

$$p^\mu := m u^\mu = \begin{pmatrix} \gamma m \\ \gamma m v^i \end{pmatrix}$$

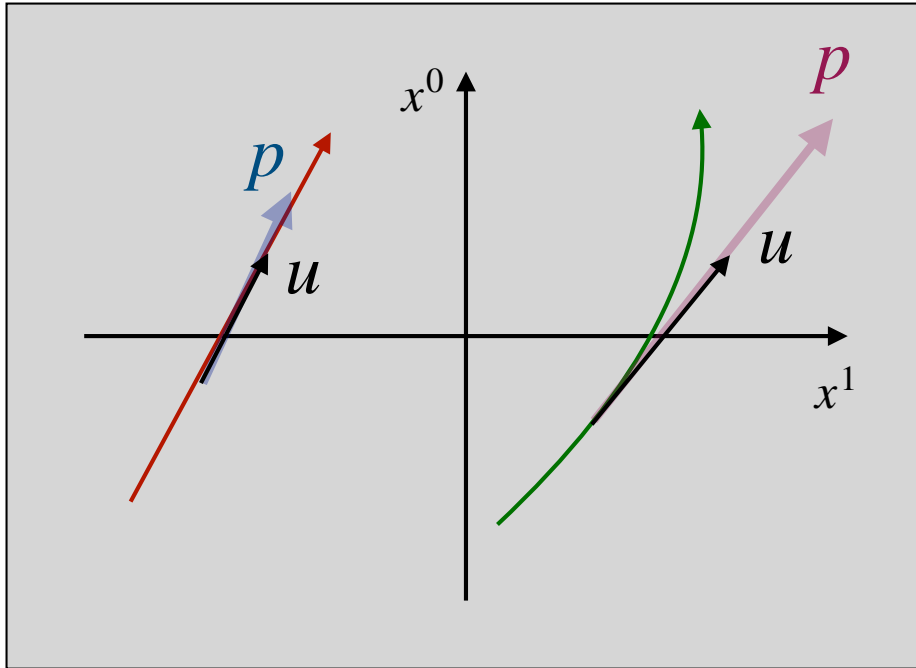
$$p \cdot p = -m^2$$

energy

momentum

total 4-momentum conserved by local forces

Kinematics in SR



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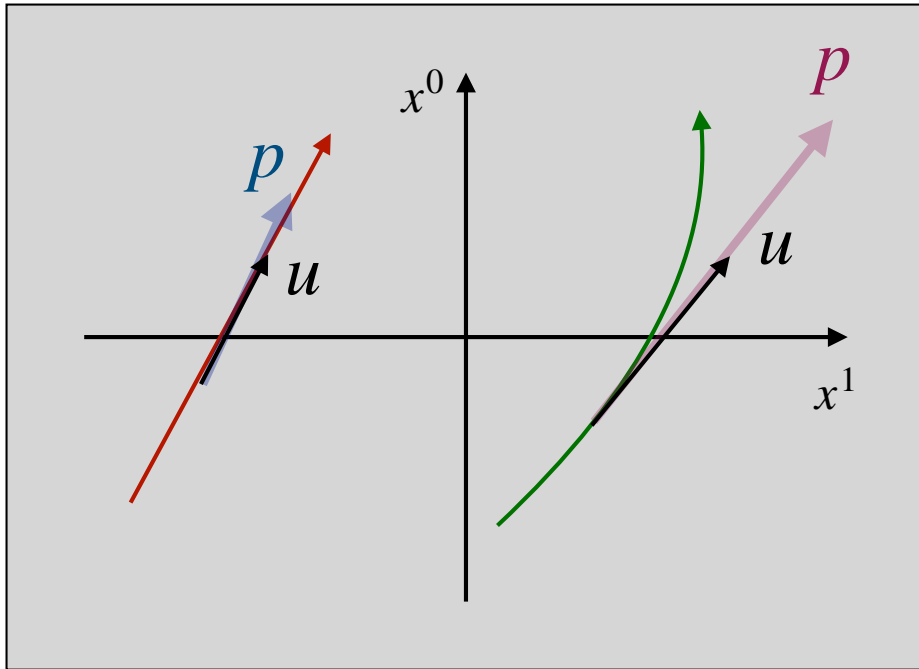
total 4-momentum conserved by local forces

small velocity limit $v \ll 1$ $v_{old}/c \ll 1$

$$\gamma = (1 - \vec{v}^2)^{-1/2} = 1 + \frac{1}{2} \vec{v}^2 + O(v^3)$$

$$p^\mu = \begin{pmatrix} m + \frac{1}{2} m \vec{v}^2 + O(v^3) \\ m v^i + O(v^3) \end{pmatrix}$$

Kinematics in SR



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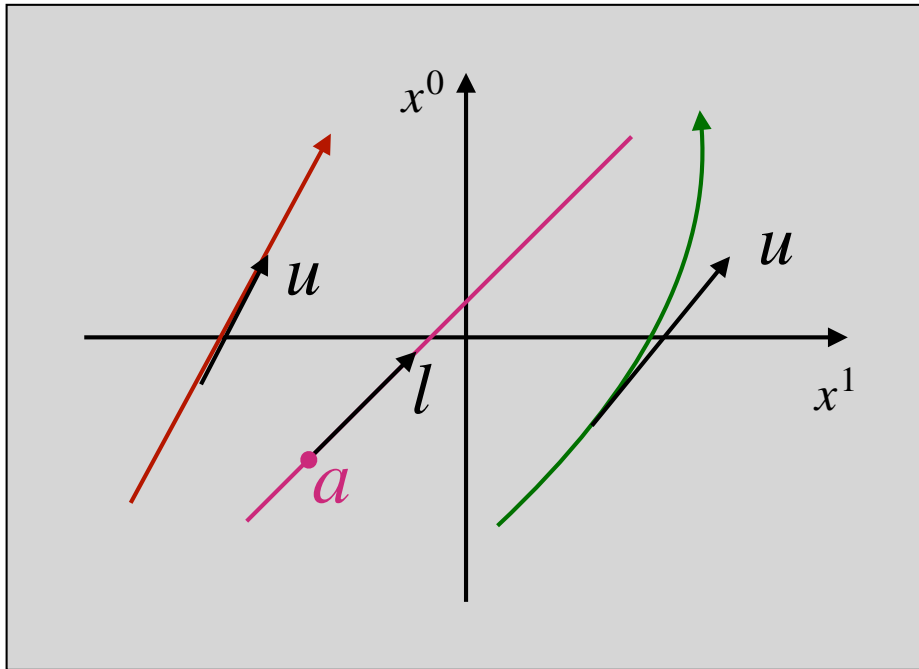
Re-introducing c

$$p_{new}^\mu = \begin{pmatrix} m + \frac{1}{2c^2} m \vec{v}_{old}^2 + O((v_{old}/c)^3) \\ m v_{old}^i / c + O((v_{old}/c)^3) \end{pmatrix}$$

$$E_{old} = c^2 p^0 = m c^2 + \frac{1}{2} m \vec{v}_{old}^2 + O((v_{old}/c)^3)$$

$$p_{old}^i = c p^i = m v_{old}^i + O((v_{old}/c)^3)$$

Kinematics in SR



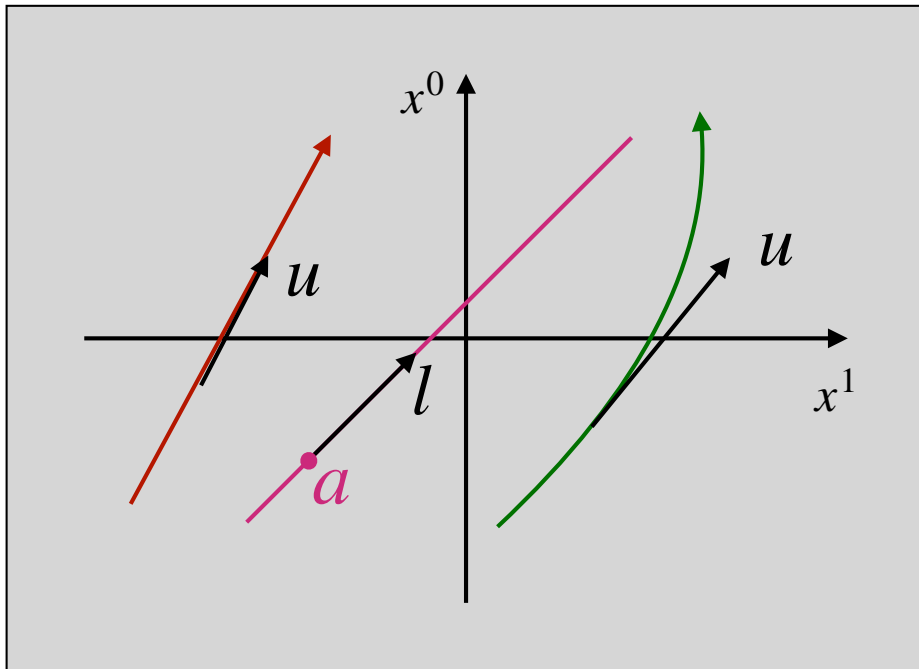
light rays (worldlines of photons)

$$x^\mu(\lambda) = a^\mu + \lambda l^\mu$$

$$l \cdot l = 0$$

$$-(l^0)^2 + \sum_i l^i l^i = 0$$

Kinematics in SR



light rays (worldlines of photons)

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velocity = speed of light, i.e. 1

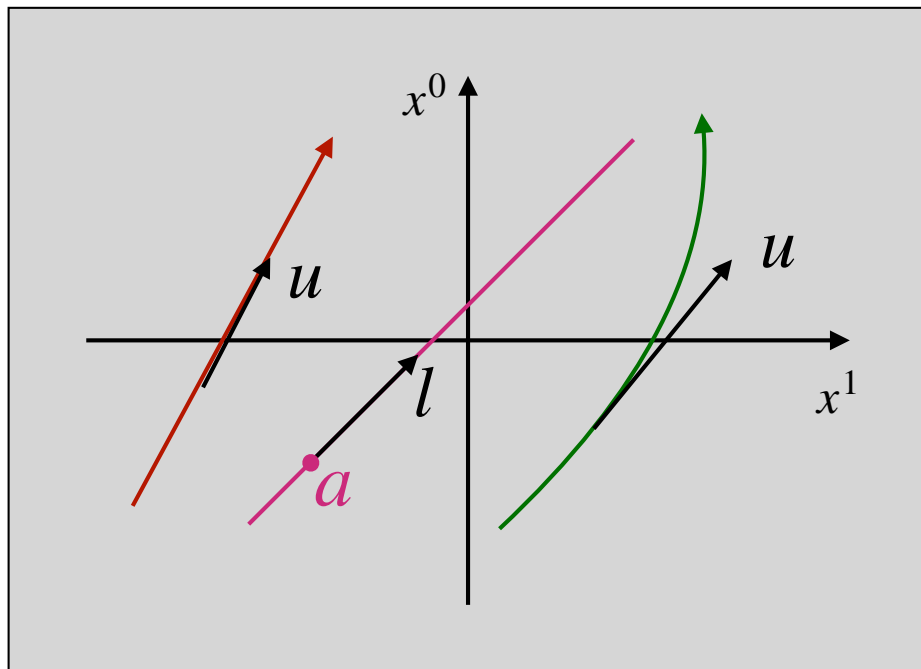
$$\frac{dx^i}{dx^0} = \frac{dx^i}{d\lambda} \cdot \frac{d\lambda}{dx^0} = \frac{l^i}{l^0}$$

$$\sum_i \frac{l^i}{l^0} \frac{l^i}{l^0} = 1$$

$$l^0 = \sqrt{\sum_i l^i l^i}$$

...in any inertial frame!

Kinematics in SR



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affine reparametrizations

$$\lambda \rightarrow \tilde{\lambda} = C\lambda + D$$

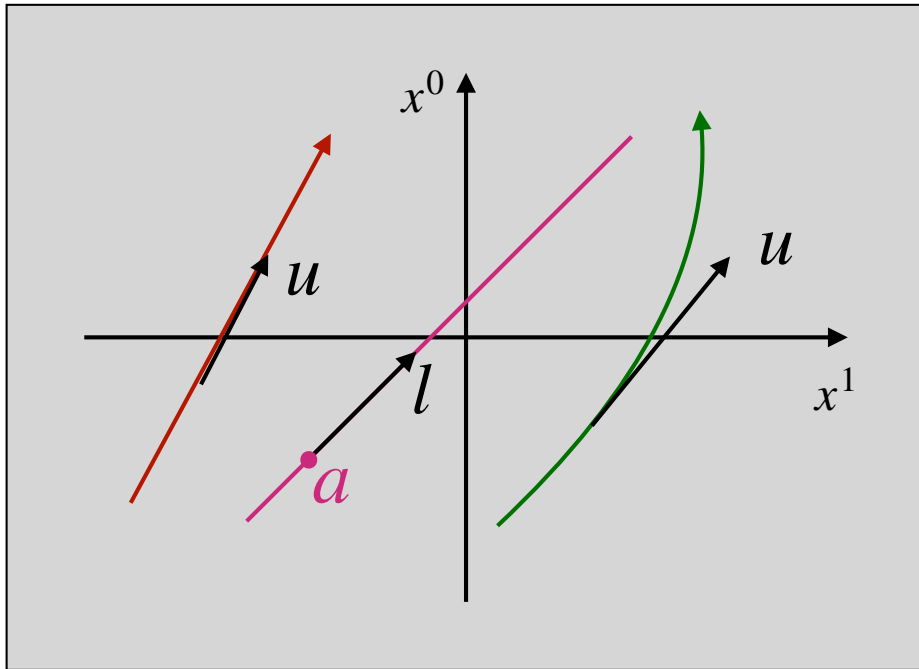
$$a \rightarrow \tilde{a} = a - D l$$

$$l \rightarrow \tilde{l} = \frac{1}{C} l$$

l not normalizable

\implies no 4-velocity or proper time for photons

Kinematics in SR



light rays (worldlines of photons)

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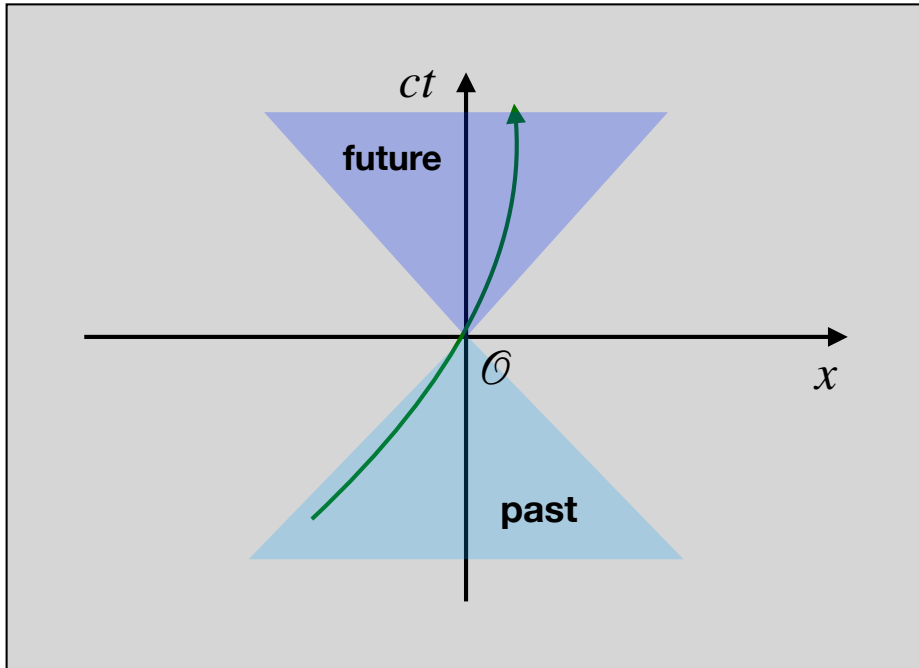
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photon's 4-momentum

$$p^\mu := \frac{E}{l^0} l^\mu \quad p \cdot p = 0$$

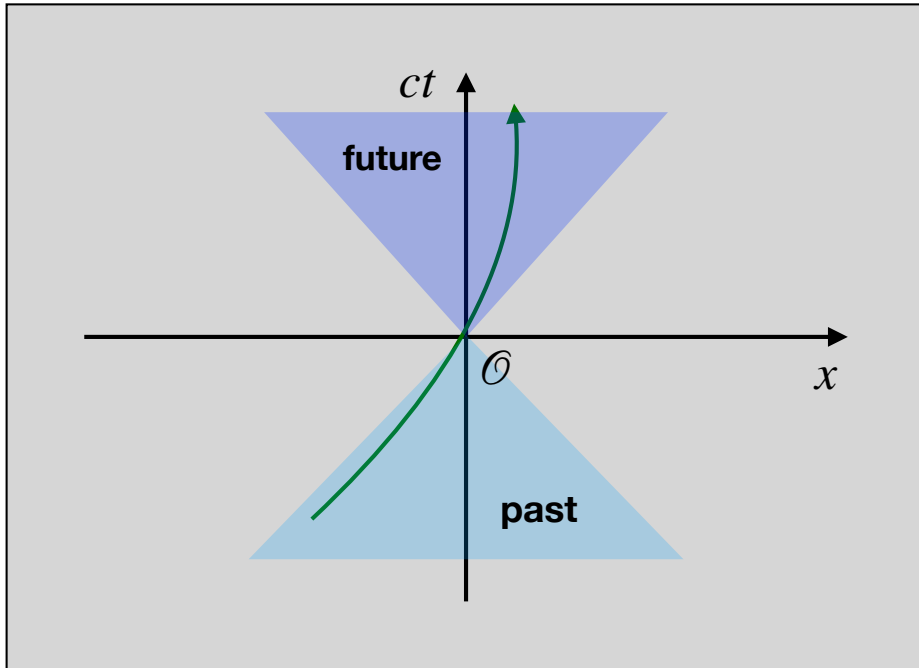
$$p^\mu = \begin{pmatrix} E \\ \vec{p} \end{pmatrix} \quad E = |\vec{p}|$$

Kinematics in SR



light cones, causality

Kinematics in SR



light cones, causality

$$\Delta x^\mu \Delta x^\nu \eta_{\mu\nu} < 0 \quad \text{timelike}$$

Inside the light cone

Worldlines of massive particles

Future/past distinction

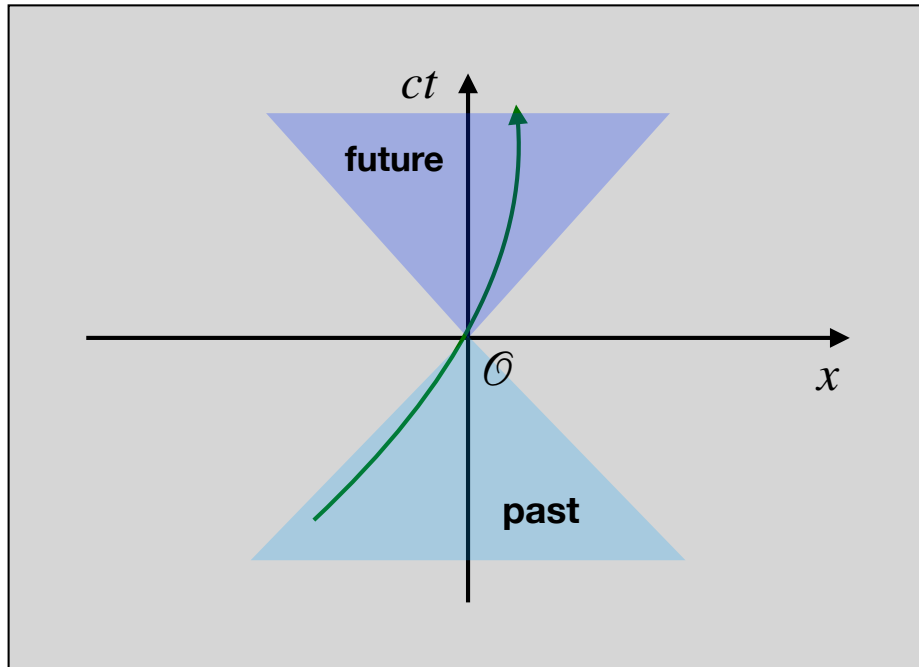
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The light cone

Points connected by light rays

Future/past distinction

Kinematics in SR



$$\Delta x^\mu \Delta x^\nu \eta_{\mu\nu} > 0 \quad \text{spacelike}$$

Outside the light cone

No future/past distinction (somewhere else)

light cones, causality

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Inside the light cone

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Future/past distinction

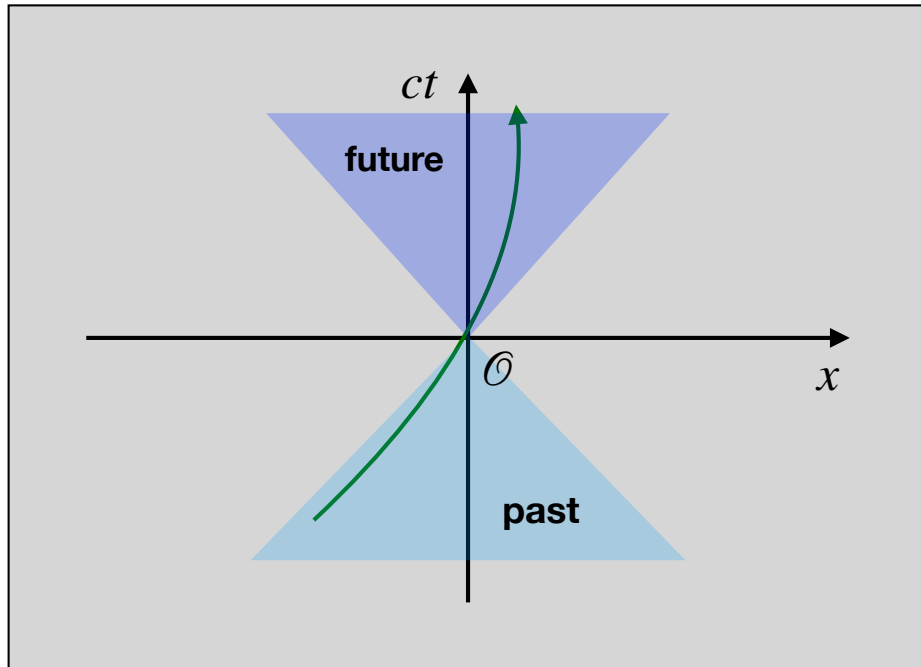
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The light cone

Points connected by light rays

Future/past distinction

$$\Delta x^\mu \Delta x^\nu \eta_{\mu\nu} > 0 \quad \text{spacelike}$$

Outside the light cone

No future/past distinction (somewhere else)

No superluminal interactions:

no worldlines outside the light cone

events outside past light cone cannot influence \mathcal{O}

\mathcal{O} cannot influence events outside future light cone

Special relativity

End of lecture 1